



**Fermilab**

TM-835

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MESON SUPERCONDUCTING MAGNET ENERGY DUMP SYSTEM

J. B. Stoffel  
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## INTRODUCTION

This paper describes an energy dump system for use with superconducting magnets when used as external beam line components.

## PROBLEM

Previous energy dump schemes needed an inductance which continuously passed the magnet current of up to 4500 amperes. In addition the dump resistor could only be placed on the power supply side of the reversing switch. This reversing switch requirement is necessary to beam line magnet systems.

## SOLUTION

The described dump system, shown in Figure I, overcomes both of these problems. An air core inductor is used and because it passes current only for the few seconds of a dump, a 1/4" copper rod conductor is sufficient. The circuit is such that the dump resistor can be placed on the magnet side of the reversing switch and thereby protect the magnet if a faulty switch opens with current flowing.

### CIRCUIT OPERATION

During operation of the superconducting magnet, the state of the system is shown in Figure II. The magnet current flows from power supply (+), through the magnet (s), through  $SCR_r$ , and back to power supply (-). The current through the dump resistor is zero as the voltage across the magnet is zero.  $C_d$  is discharged, and  $C_c$  is charged with  $550^V (V_{cc})$

When  $SCR_c$  is fired, this charge will reverse bias  $SCR_r$  for the required  $250\mu$  sec so that it will turn off reliably and force the magnet current through the dump resistor  $R_d$ . The purpose of  $C_d$  is to limit the  $dv/dt$  that  $SCR_r$  sees as the current is diverted through  $R_d$ .  $L_d$  limits the current through  $R_d$  during the reverse bias of  $SCR_r$ . The short and long term waveforms are shown in Figure IX.

### DETAILED ANALYSIS

Figure III shows our system model and its state before the dump sequence starts. The power supply is driving current through the magnet via the run SCR's ( $SCR_r$ ). A closer look (figure I) shows six (6) SCR's, in parallel, each with sharing resistors.  $C_c$  is charged with a voltage  $V_{cc}$ ,  $C_d$  is discharged, and  $i_{c1} = 0$  in the commutation loop. Also, note that  $i_d = 0$  in  $L_d$  (dump inductance) and  $R_d$  (dump resistor) as the voltage across the superconducting magnet =  $0^V$ .

The dump is initiated when the gate pulses are removed from  $SCR_r$  and  $SCR_c$  is fired. Figure IV shows the model used to study the commutation loop current buildup. Initially  $i_m$  is flowing thru  $SCR_r$ . Then with a slope of  $(V_{cc} - V_{cd}) \div L_c$ ,  $i_{c1}$  increases to  $i_m$ .  $L_c$  is included in the loop to prevent abrupt changes in  $i_{c1}$  ( $\frac{di}{dt} < 100A/\mu s$ ). At first  $V_{cd} = 0$ , but as  $i_{c1}$  continues to flow  $V_{cc}$  decreases as  $V_{cd}$  increases with the polarity shown in Figure IV. When  $i_{c1} = i_m$ , SCR becomes reversed biased and another model must be considered.

Figure VII is a graphical display of loop current and capacitor voltage for different values of loop inductance and capacitance. With a commutation voltage of  $550^V$ , an  $L_c = 8\mu h$  limits  $\frac{di}{dt}$  thru  $SCR_c$  to  $69 \text{ amps}/\mu s$ . This is 69% of the maximum condition to achieve long industrial life ( $\geq 20$  years).

Figure V is the model used to analyze the system for  $SCR_r$  reversed biased. The analysis is straightforward if the magnet is modeled as a current source with  $i = i_m$ . This is valid for times that are small compared to  $L_m/R_d$ .

Figure IX shows the voltages and currents of interest that are generated by the model when the actual component values are used. At the instant  $SCR_r$  becomes reversed biased, the voltage across it jumps to a maximum and then decays as a function of the capacitance and the magnet current. As  $C_c$  discharges,  $C_d$  charges and  $SCR_r$  becomes forward biased. If the reverse bias period was more than 250  $\mu$  sec it won't conduct. The  $C_d$  charges up and the current through the dump resistor increases to its maximum value of  $i_m$ . Note that this happens in 5 m sec which is relatively small compared to the magnet-dump resistor time constant (.675 sec).

Figure VIII shows the effect that varying the capacitance will have on the time of reverse bias for  $SCR_r$ . To guarantee commutation, this time must be in excess of 250  $\mu$ s. Hence, for  $i_m = 4500$  amp,  $L_c = 8 \mu$ h,  $C_c = 16,000 \mu$ f, and  $C_d = 6400 \mu$ f the duration of the reverse bias will be a conservative 383  $\mu$  sec.

The simple model shown in Figure VI is adequate for a long term study on the dump voltage as all other system time constants are relatively short. Figure IX shows this dump voltage

# EXPERIMENTAL RESULTS

The pictures in Figure X of the  $SCR_r$  and  $R_d$  voltages during a dump sequence shows a good approximation to Figure IX and its calculated waveforms.

The table below demonstrates that the predicted values of  $\Delta T_{rb}$  (reverse bias) are longer than the measured ones.

<u>I Magnet</u>	<u><math>\Delta T</math> (calculated)</u>	<u><math>\Delta T</math> (measured)</u>
2000 A	750 $\mu s$	720 $\mu s$
2500 A	650 $\mu s$	600 $\mu s$
3000 A	560 $\mu s$	500 $\mu s$
3500 A	500 $\mu s$	440 $\mu s$

The top picture also shows an overshoot when the SCR is reverse biased. This can be explained by  $i_{c1}$  increasing to more than  $i_m(0)$  due to the charge carriers that must be swept out of the junctions of  $SCR_r$  during the process of reversing the bias.

Figure XI shows pictures of the dump voltage with and without a magnet quench. Note the fast decay of the quench photo as compared to a non-quench. This is due to the increasing resistance in the magnet as the quench propagates.

The author received a great deal of assistance in this project from:

John Dinkel explained the operation of the original dump scheme and was a source of parts to get the project off to a fast start.

Paul Czarapata helped with his knowledge of programming the PDP-10 for the analysis of the system.

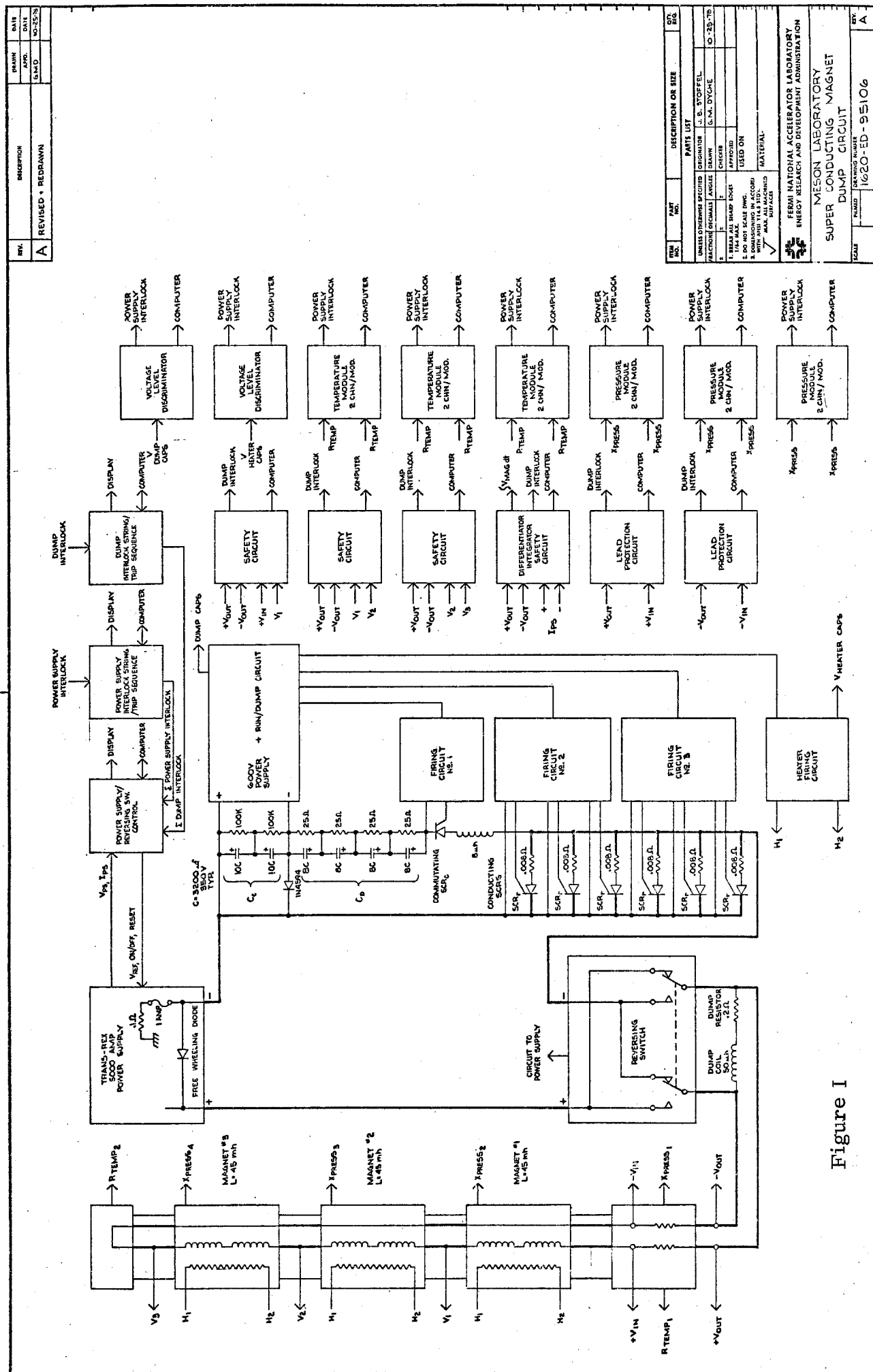
Terrance O'Brien made the greatest contribution by building and debugging the system.

And in addition many other people at the Meson Lab, Energy Doubler and Accelerator.

References: (For the detailed analysis)

1. Charles M. Close  
"The Analysis of Linear Circuits"; Chapter 10  
Harcourt Brace and World Inc. 1966.
2. Sylvan Fich  
"Transient Analysis in Electrical Engineering"  
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McGraw-Hill Book Company, 1975





## Figure I



# SIMPLIFIED SYSTEM OVERVIEW

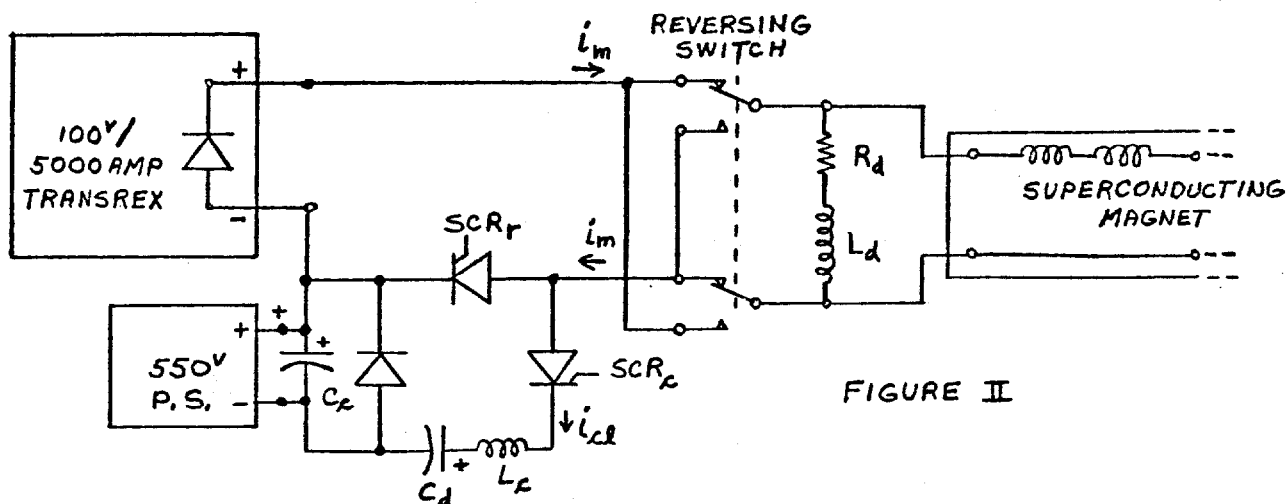


FIGURE II

$R_d$  = DUMP RESISTOR       $i_m$  = MAGNET CURRENT       $SCR_c$  = COMMUTATION SCR  
 $L_d$  = DUMP INDUCTANCE       $i_{cl}$  = COMMUTATION LOOP CURRENT       $SCR_r$  = CONDUCTION SCR  
 $C_c$  = COMMUTATION CAP.       $C_d$  = DUMP CAPACITOR       $L_c$  = COMMUTATION INDUCTANCE

## CONDUCTION MODEL

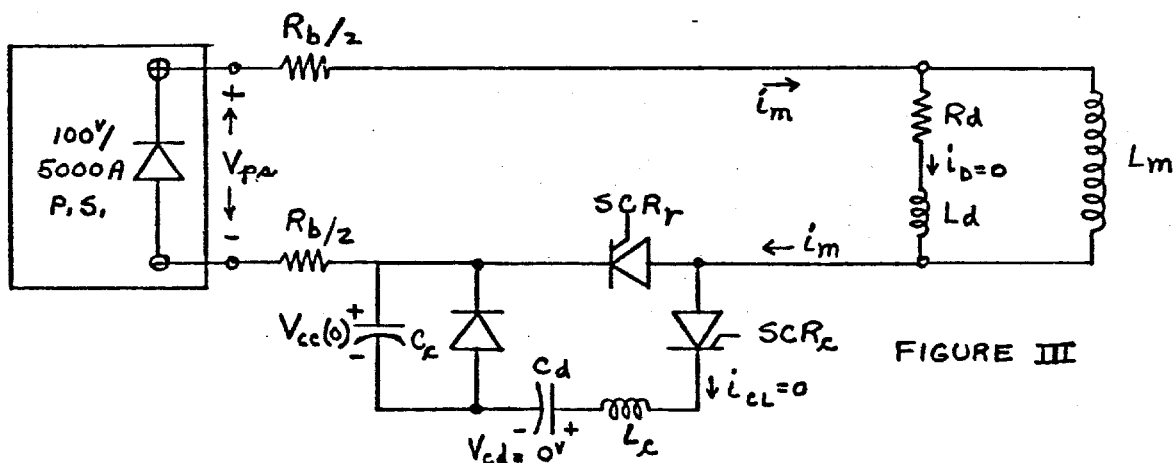


FIGURE III

$L_m$  = MAGNET INDUCTANCE       $R_b$  = BUSS RESISTANCE       $i_d$  = CURRENT THRU  $R_d$   
 $V_{p\Delta}$  = POWER SUPPLY VOLTAGE       $V_{cc}(0)$  = INITIAL VOLTAGE ON  $C_c$        $V_{cd} = 0V$



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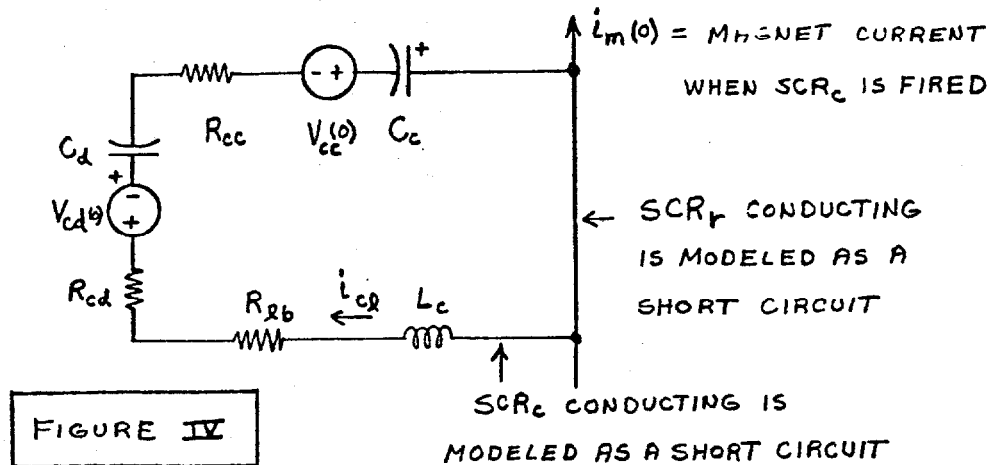
FIGURE IV

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COMMUTATION LOOP CURRENT BUILDUP MODEL



$R_{cc} = \text{RESISTANCE OF } C_c$   $R_{cd} = \text{RESISTANCE OF } C_d$

$R_{lb} = \text{RESISTANCE OF LOOP BUSS}$

$V_{cc}(0) = \text{INITIAL VOLTAGE ON } C_c$   $V_{cd}(0) = \text{INITIAL VOLTAGE ON } C_d$

FOR THE CURRENT BUILDUP MODEL:

$$i_{cl}(t) = \frac{(V_{cc}(0) - V_{cd}(0))}{\omega L_c} e^{-\alpha t} \sin \omega t$$

$$V_{cc}(t) = V_{cc}(0) - \frac{(V_{cc}(0) - V_{cd}(0))}{C_c \omega} \sqrt{\frac{L_c}{C}} \sin(\omega t + \beta) - L_c C$$

$$V_{cd}(t) = V_{cd}(0) + \frac{(V_{cc}(0) - V_{cd}(0))}{C_d \omega} \sqrt{\frac{L_c}{C}} \sin(\omega t + \beta) + L_c C$$

WHERE  $C = \frac{C_c \times C_d}{C_c + C_d}$   $R = R_{cc} + R_{cd} + R_{lb}$

$$\alpha = \frac{R}{2L_c} \quad \omega = \left( \frac{1}{L_c C} - \frac{R^2}{4L_c^2} \right)^{1/2} \quad \beta = \tan^{-1} \frac{\omega}{\alpha} - \pi$$

AS THE UNDERDAMPED CASE  $\left( \frac{1}{L_c C} > \frac{R^2}{4L_c^2} \right)$  IS THE ONE OF PRACTICAL INTEREST



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FIGURES V AND VI

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SCR<sub>R</sub> REVERSED BIASED MODEL

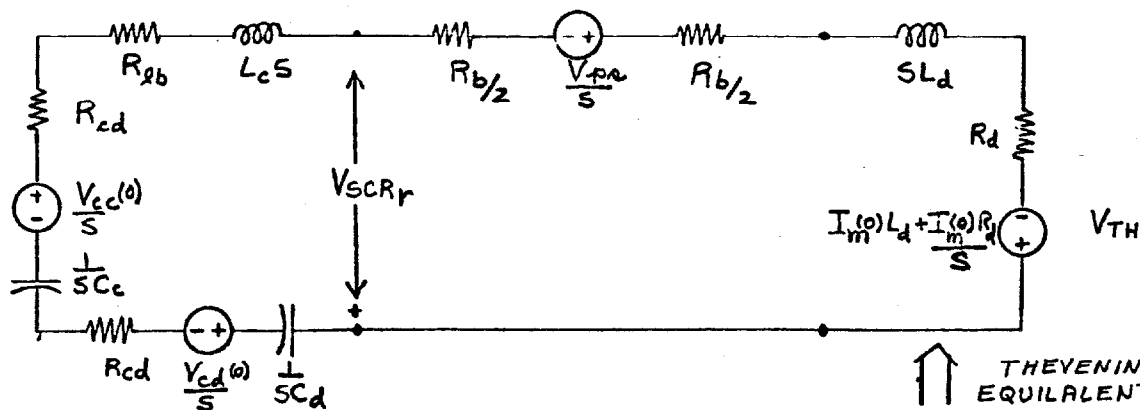
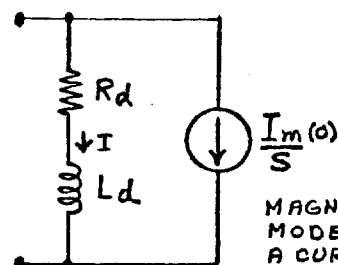


FIGURE V

MODEL EQUATIONS ON NEXT PAGE

THEVENIN'S  
EQUIVALENT CKT.



MAGNETS ARE  
MODELLED AS  
A CURRENT SOURCE

LONG TERM DUMP MODEL

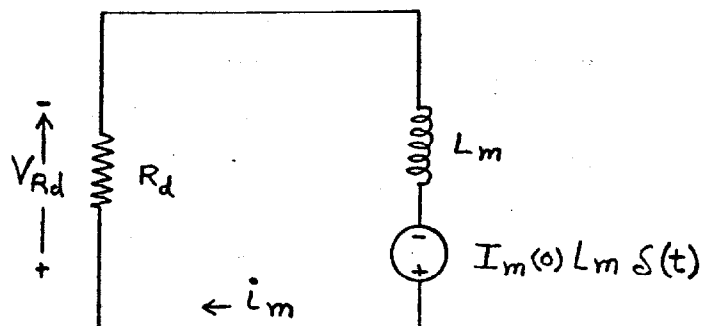


FIGURE VI

$$i_m(t) = I_m(0) e^{-t \frac{R_d}{L_d}}$$

$$V_{R_d}(t) = R_d i_m(t)$$



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FIGURE VII

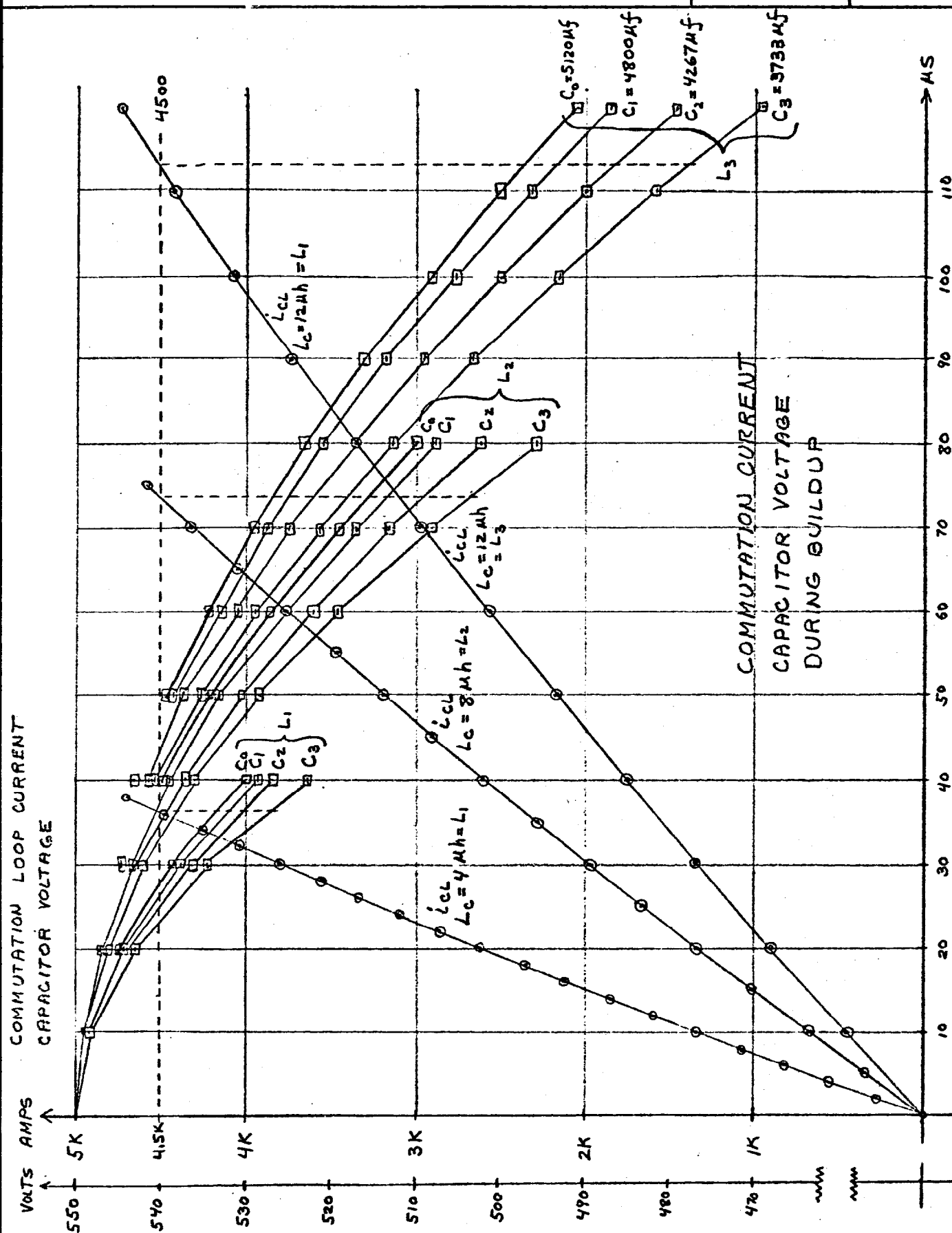


FIGURE VII



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FIGURE VIII

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## REVERSE BIAS STUDIES

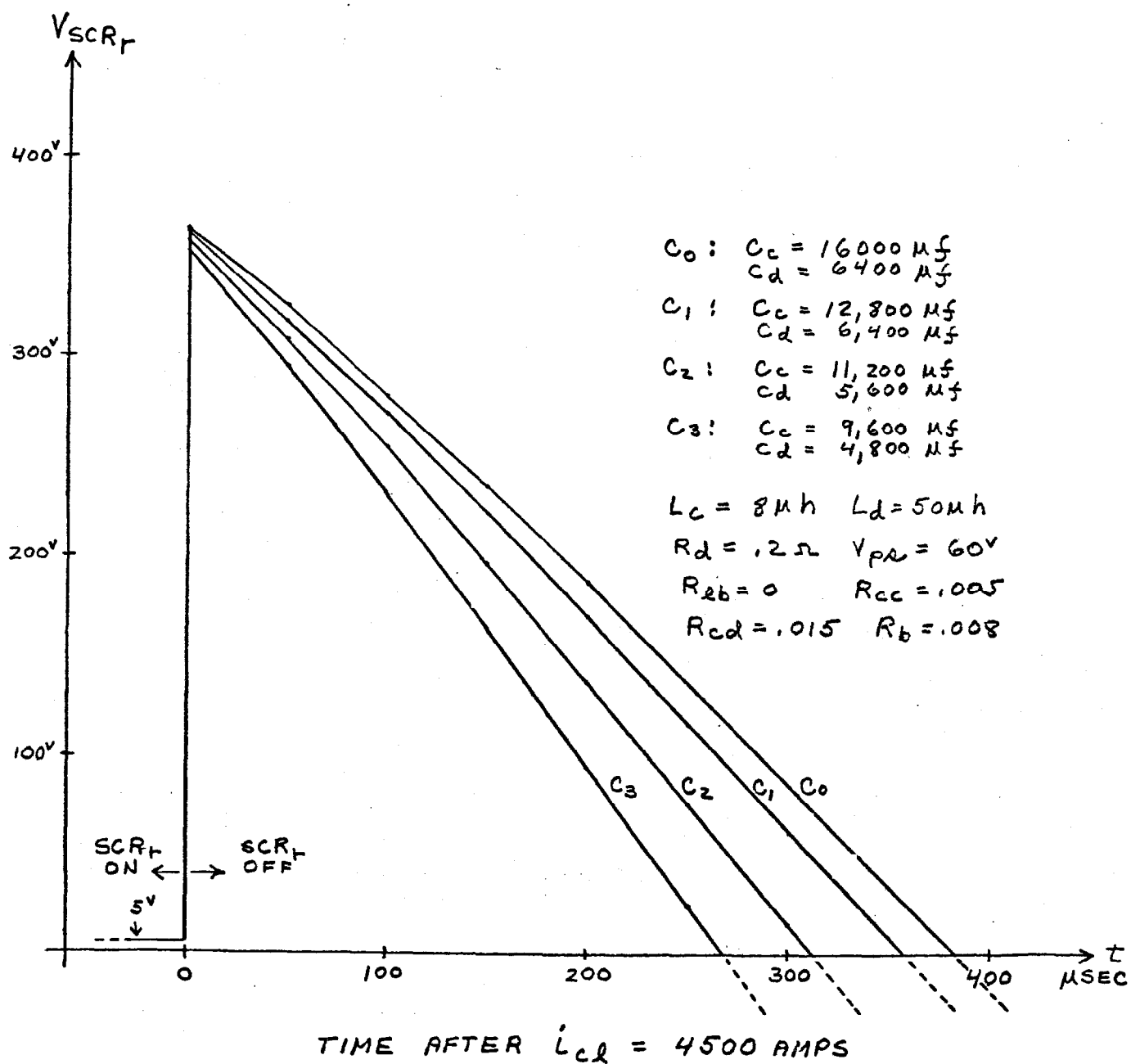


FIGURE VIII



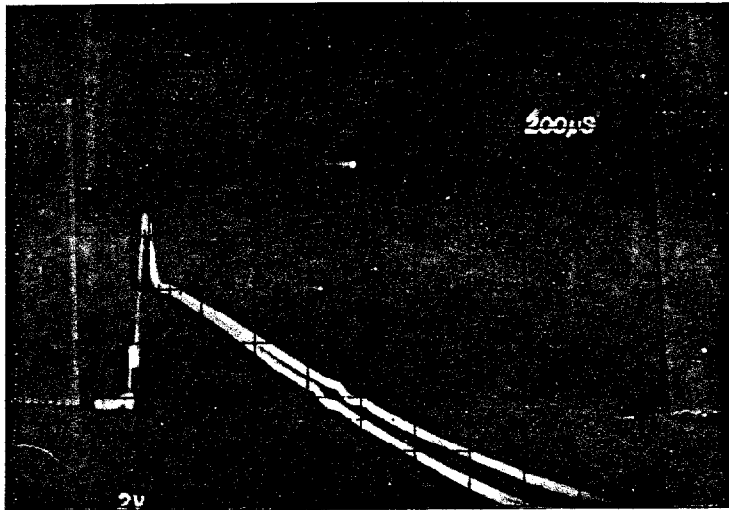
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FIGURE X

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UPPER  $I_m(0) = 2000$  AMPS

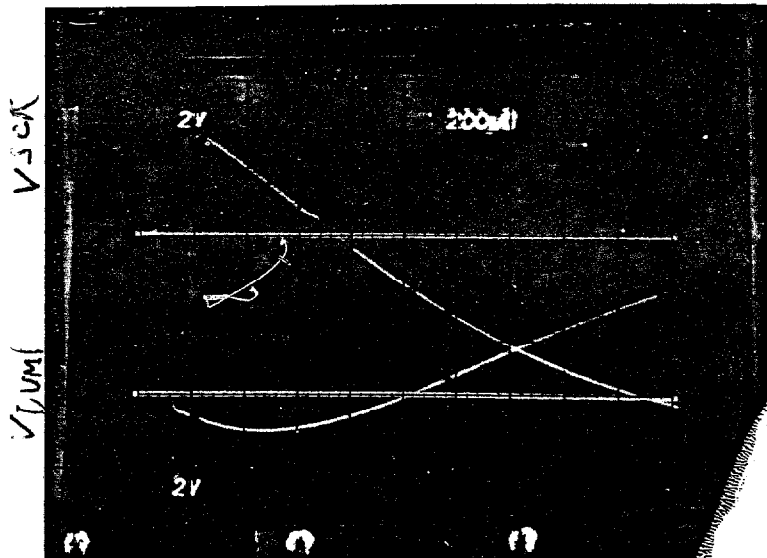
LOWER  $I_m(0) = 2500$  AMPS

$V_{SCRr}$  (REV BIAS)  
200V/DIV

200µS/DIV

$\Delta T(RB) = 720$  MS @  $I_m(0) = 2000$  A

$\Delta T(RB) = 600$  MS @  $I_m(0) = 2500$  A



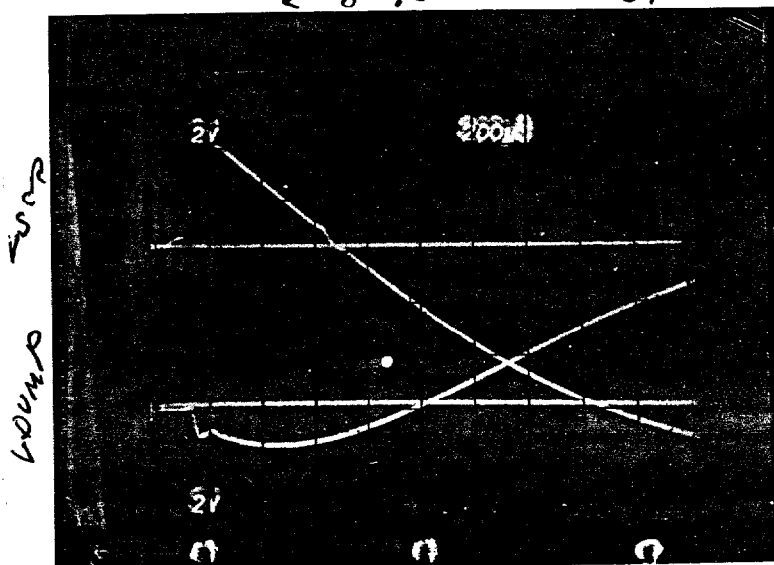
$V_{SCRr}$  200V/DIV

200µS/DIV

$V_{Rd}$  200V/DIV

200µS/DIV

$\Delta T(RB) = 500$  MS @  $I_m = 3000$  A



$V_{SCRr}$  200V/DIV

200µS/DIV

$V_{Rd}$  200V/DIV

200µS/DIV

$\Delta T(RB) = 440$  MS @ 3500 A



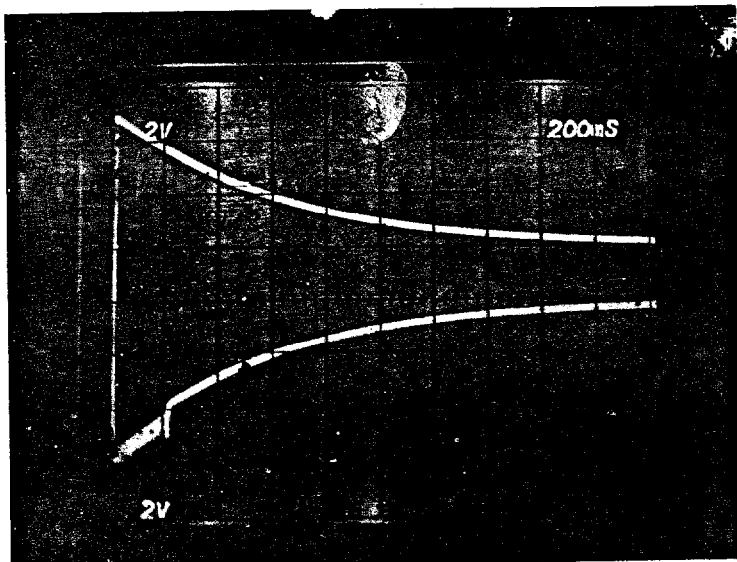
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FIGURE XI

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MAGNET NOT QUENCHED

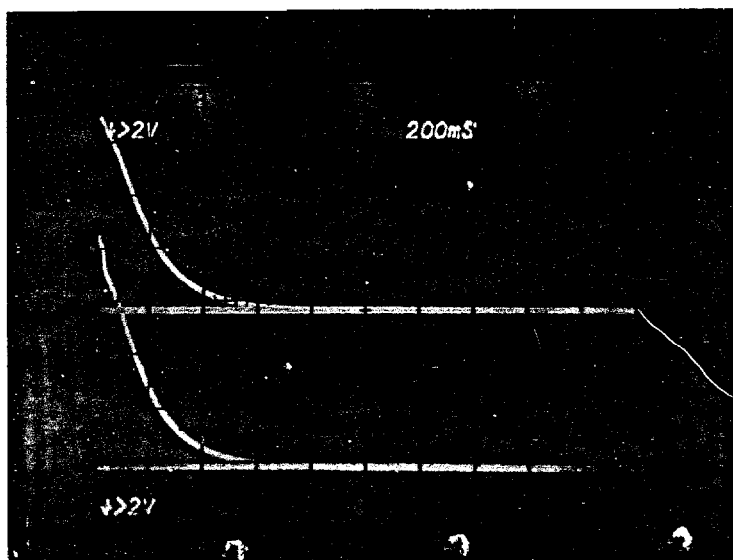
6/30/78

$I_m(0) = 2000 A$   
 $V_{R_d} (200V/DIV)$   
 $\rightarrow 200 ms / DIV$   
 $V_{SCR_F} (FORWARD BIAS) (200V/DIV)$   
 $\rightarrow 200 ms / DIV$

MAGNET QUENCHED

$I_m(0) = 3500 A$

10/25/78



$V_{R_d} (200V/DIV)$   
 $\rightarrow 200 ms / DIV$   
 $V_{SCR_F} (200V/DIV)$   
 $\rightarrow 200 ms / DIV$